

MIT OpenCourseWare
<http://ocw.mit.edu>

5.04 Principles of Inorganic Chemistry II
Fall 2008

For information about citing these materials or our Terms of Use, visit: <http://ocw.mit.edu/terms>.

With these parameters defined, the polynomial form of $D_N(x)$ for any value of N can be obtained,

$$D_3(x) = xD_2(x) - D_1(x) = x(x^2-1) - x = x(x^2-2)$$

$$D_4(x) = xD_3(x) - D_2(x) = x^2(x^2-2) - (x^2-1)$$

⋮

and so on

The expansion of $D_N(x)$ has as its solution,

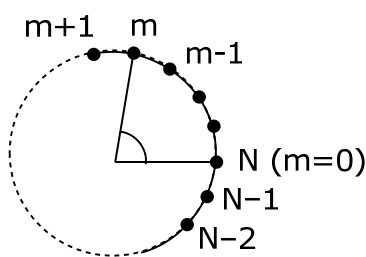
$$x = -2 \cos \frac{2\pi}{N} j \quad (j = 0, 1, 2, 3 \dots N-1)$$

and substituting for x ,

$$E = \alpha + 2\beta \cos \frac{2\pi}{N} j \quad (j = 0, 1, 2, 3 \dots N-1)$$

Standing Wave Derivation

An alternative approach to solving this problem is to express the wavefunction directly in an angular coordinate, θ



For a standing wave of λ about the perimeter of a circle of circumference c ,

$$\psi_j = \sin \frac{c}{\lambda} \theta$$

The solution to the wave function must be single valued \therefore a single solution must be obtained for ψ at every $2n\pi$ or in analytical terms,

$$\begin{aligned} \psi &= \sin \frac{c}{\lambda} (\theta + 2\pi) = \sin \frac{c}{\lambda} \theta \\ &= \sin \frac{c}{\lambda} \theta \cdot \cos \frac{c}{\lambda} 2\pi + \sin \frac{c}{\lambda} 2\pi \cdot \cos \frac{c}{\lambda} \theta = \sin \frac{c}{\lambda} \theta \end{aligned}$$

\downarrow must go to 1 \downarrow must go to 0

iff $\frac{c}{\lambda} 2\pi = 2\pi j$ ($j = 0, 1, 2 \dots N-1$)

$\therefore \frac{c}{\lambda} = j$ ↗ condition for an integral number of λ 's about the circumference of a circle

Thus the amplitude of ψ_j at atom m is, (where $\frac{c}{\lambda} = j$ and $\theta = \frac{2\pi}{N}m$)

$$\psi_j(m) = \sin \frac{2\pi m}{N} j \quad (j = 0, 1, 2 \dots N-1)$$

Within the context of the LCAO method, ψ_j may be rewritten as a linear combination in ϕ_m with coefficients c_{jm} . Thus the amplitude of ψ_j at m is equivalent to the coefficient of ϕ_m in the LCAO expansion,

$$\begin{aligned} \psi_j &= \sum_{m=1}^N c_{jm} \phi_m \\ \text{where } c_{jm} &= \sin \frac{2\pi m}{N} j \quad (j = 0, 1, 2 \dots N-1) \end{aligned}$$

The energy of each MO, ψ_j , may be determined from a solution of Schrödinger's equation,

$$\begin{aligned} H\psi_j &= E_j\psi_j \\ \left| H - E_j \right| \psi_j \rangle &= 0 \\ \left| H - E_j \right| \sum_m^N c_{jm} \phi_m \rangle &= 0 \end{aligned}$$

The energy of the ϕ_m orbital is obtained by left-multiplying by ϕ_m ,

$$\langle \phi_m | H - E_j | \sum_m^N c_{jm} \phi_m \rangle = 0$$

but the Hückel condition is imposed; the only terms that are retained are those involving ϕ_m , ϕ_{m+1} , and ϕ_{m-1} . Expanding,

$$\begin{aligned} & \left[c_{jm} \langle \phi_m | \overset{\alpha}{H} | \phi_m \rangle - c_{jm} E_j \langle \phi_m | \overset{1}{\phi_m} \rangle \right] + \left[c_{j(m+1)} \langle \phi_m | \overset{\beta}{H} | \phi_{m+1} \rangle - c_{j(m+1)} E_j \langle \phi_m | \overset{0}{\phi_{m+1}} \rangle \right] \\ & + \left[c_{j(m-1)} \langle \phi_m | \overset{\beta}{H} | \phi_{m-1} \rangle - c_{j(m-1)} E_j \langle \phi_m | \overset{0}{\phi_{m-1}} \rangle \right] = 0 \end{aligned}$$

Evaluating the integrals,

$$\alpha c_{jm} - c_{jm} E_j + \beta [c_{j(m+1)} + c_{j(m-1)}] = 0$$

$$\alpha c_{jm} + \beta [c_{j(m+1)} + c_{j(m-1)}] = c_{jm} E_j$$

Substituting for c_{jm} ,

$$\alpha \sin \frac{2\pi m}{N} j + \beta \left(\sin \frac{2\pi(m+1)}{N} j + \sin \frac{2\pi(m-1)}{N} j \right) = E_j \sin \frac{2\pi m}{N} j$$

Dividing by $\sin \frac{2\pi m}{N} j$,

$$\alpha + \frac{\beta \left(\sin \frac{2\pi(m+1)}{N} j + \sin \frac{2\pi(m-1)}{N} j \right)}{\sin \frac{2\pi m}{N} j} = E_j$$

Making the simplifying substitution, $\kappa = \frac{2\pi}{N} j$

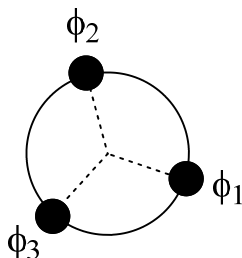
$$E_j = \alpha + \frac{\beta (\sin \kappa(m+1) + \sin \kappa(m-1))}{\sin \kappa m}$$

$$E_j = \alpha + \beta \left(\frac{\sin \kappa m \cdot \cos \kappa + \sin \kappa \cdot \cos \kappa m + \sin \kappa m \cdot \cos \kappa - \sin \kappa \cdot \cos \kappa m}{\sin \kappa m} \right)$$

$$E_j = \alpha + 2\beta \cos \kappa$$

$$E_j = \alpha + 2\beta \cos \frac{2\pi}{N} j \quad (j = 0, 1, 2 \dots N-1)$$

Let's look at the simplest cyclic system, $N = 3$



$$N = 3, \text{ so } E_j = \alpha + 2\beta \cos \frac{2\pi}{N} j \text{ where } j = 0, 1, 2$$

$$E_0 = \alpha + 2\beta$$

$$E_1 = \alpha + 2\beta \cos \frac{2\pi}{3} = \alpha - \beta$$

$$E_2 = \alpha + 2\beta \cos \frac{4\pi}{3} = \alpha - \beta$$

Continuing with our approach (LCAO) and using E_j to solve for the eigenfunction, we find...

$$\psi_j = \sum_m e^{ij\theta} \phi_m \quad \text{for } j = 0, \pm 1, \pm 2 \dots \begin{cases} \pm \frac{N}{2} \text{ for } N \text{ even} \\ \pm \frac{(N-1)}{2} \text{ for } N \text{ odd} \end{cases}$$

Using the general expression for ψ_j , the eigenfunctions are:

$$\psi_0 = e^{i(0)0} \phi_1 + e^{i(0)\frac{2\pi}{3}} \phi_2 + e^{i(0)\frac{4\pi}{3}} \phi_3$$

$$\psi_{+1} = e^{i(1)0} \phi_1 + e^{i(1)\frac{2\pi}{3}} \phi_2 + e^{i(1)\frac{4\pi}{3}} \phi_3$$

$$\psi_{-1} = e^{i(-1)0} \phi_1 + e^{i(-1)\frac{2\pi}{3}} \phi_2 + e^{i(-1)\frac{4\pi}{3}} \phi_3$$

Obtaining real components of the wavefunctions and normalizing,

$$\psi_0 = \phi_1 + \phi_2 + \phi_3 \rightarrow \psi_0 = \frac{1}{\sqrt{3}} (\phi_1 + \phi_2 + \phi_3)$$

$$\psi_{+1} + \psi_{-1} = 2\phi_1 - \phi_2 - \phi_3 \rightarrow \psi_1 = \frac{1}{\sqrt{6}} (2\phi_1 - \phi_2 - \phi_3)$$

$$\psi_{+1} - \psi_{-1} = \phi_2 - \phi_3 \rightarrow \psi_2 = \frac{1}{\sqrt{2}} (\phi_2 - \phi_3)$$

Summarizing on a MO diagram where α is set equal to 0,

