

# Accuracy & Consistency for Finite Difference Methods

We begin with defining a convergent method:

Definition: A finite difference method for solving  
 Initial Value Problem  $\Rightarrow \frac{d\vec{u}}{dt} = \vec{F}(\vec{u}, t) \quad \vec{u}(0) = \vec{u}_0$  } We will call this our IVP from now on

from  $t=0$  to  $T$  is convergent if

$$\max_{n \in [0, \frac{T}{\Delta t}]} \|\vec{v}^n - \vec{u}(n\Delta t)\| \rightarrow 0 \quad \text{as } \Delta t \rightarrow 0$$

The global order of accuracy measures how quickly the finite difference method converges:

Definition: The global order of accuracy for a finite difference method for IVP is  $p$  when:

$$\max_{n \in [0, \frac{T}{\Delta t}]} \|\vec{v}^n - \vec{u}(n\Delta t)\| \leq C (\Delta t)^p \quad \text{as } \Delta t \rightarrow 0$$

Assumes  $\vec{F}$  has continuous derivatives up to and including  $\frac{\partial^p \vec{F}}{\partial t^p}$  and  $\frac{\partial^p \vec{F}}{\partial \vec{u}^p}$ .

Another definition we require is consistency 2

Definition: A finite difference method is consistent if

$$\|\vec{v}' - \vec{v}(\Delta t)\| \leq C \Delta t \quad \text{as } \Delta t \rightarrow 0$$

for any initial condition.

Consistency is a bare minimum condition required of any numerical algorithm. If a method were not consistent, then  $\vec{v}'$  can never approach  $\vec{v}(\Delta t)$  no matter how small  $\Delta t$  is chosen.

Consistency is checked through a Taylor series analysis. Here's an example for Forward Euler:

Forward Euler:  $\vec{v}' = \vec{v}^0 + \Delta t \vec{F}(\vec{v}^0, 0)$

$$\vec{v}' - \vec{v}(\Delta t) = \vec{v}^0 + \Delta t \vec{F}(\vec{v}^0, 0) - \vec{v}(\Delta t)$$

But, we assume  $\vec{v}^0 = \vec{v}(0)$

$$\vec{v}' - \vec{v}(\Delta t) = \vec{v}(0) + \Delta t \vec{F}(\vec{v}(0), 0) - \vec{v}(\Delta t)$$

and  $\vec{F}(\vec{v}(0), 0) = \dot{\vec{v}}(0)$

$$\vec{v}' - \vec{v}(t) = \vec{v}(0) + \Delta t \dot{\vec{v}}(0) - \vec{v}(t)$$

But, a Taylor series expansion of  $\vec{v}(t)$  gives:

$$\vec{v}(t) = \vec{v}(0) + \Delta t \dot{\vec{v}}(0) + \frac{1}{2} \Delta t^2 \ddot{\vec{v}}(0) + O(\Delta t^3)$$

Thus, substituting this Taylor series gives:

$$\vec{v}' - \vec{v}(t) = -\frac{1}{2} \Delta t^2 \ddot{\vec{v}}(0) + O(\Delta t^3)$$

Clearly, this is faster than  $\Delta t$  as required by consistency.

$\Rightarrow$  P.E. is consistent

The global error can be viewed as a sum of local errors at every iteration.

Define:  $\vec{E}^n \equiv$  Global error at  $t^n \equiv \vec{v}^n - \vec{v}(n\Delta t)$

The difference between the global error  $\vec{E}^{n+1} - \vec{E}^n$  is the local error:

$$\Delta \vec{E}^n \equiv \vec{E}^{n+1} - \vec{E}^n = \vec{v}^{n+1} - \vec{v}^n - \vec{v}^{n+1} + \vec{v}^n$$

We will write  $\vec{v}^{n+1} = N(\vec{v}^n)$  where  $N(\vec{v}^n)$  is the numerical integration method:

$$\Delta \vec{E}^n(\vec{v}^n) = N(\vec{v}^n) - \vec{v}^n - \vec{v}^{n+1} + \vec{v}^n$$

$$\Rightarrow \vec{E}^n = \sum_{i=0}^n \Delta \vec{E}^i$$

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Now, suppose that  $\Delta \vec{E}^i = O(\Delta t^{p+1})$

$$\Rightarrow \vec{E}^n = O(n \Delta t^{p+1})$$

The largest error occurs for the largest  $n$  which is  $n = N = \frac{T_{\max}}{\Delta t}$

$$\Rightarrow \vec{E}^N = O\left(\frac{T_{\max}}{\Delta t} \Delta t^{p+1}\right)$$

$$\Rightarrow \boxed{\text{Global error} = O(\Delta t^p) \text{ if local error is } O(\Delta t^{p+1})}$$

Calculation of the order of  $\Delta \vec{E}^n$  is difficult. Instead, we define the closely related local truncation error (LTE) as the difference of  $N(\vec{u}^n)$  and  $\vec{u}^{n+1}$ :

$$\vec{e}^n \equiv N(\vec{u}^n) - \vec{u}^{n+1} \leftarrow \text{LTE}$$

Note:  $\vec{e}^n = \Delta \vec{E}^n(\vec{u}^n)$ .

Definition: The local order of accuracy of the numerical integration method is  $p$  where:

$$\|N(\vec{u}^n) - \vec{u}^{n+1}\| \leq C \Delta t^{p+1} \quad \text{as } \Delta t \rightarrow 0$$

and  $\vec{u} = \text{exact solution of ODE's}$ .

We note that the local order of accuracy is closely tied to the scheme consistency.

Local order of accuracy of Forward Euler

$$\vec{v}^{n+1} = \vec{v}^n + \Delta t \vec{F}(\vec{v}^n, t^n) = N(\vec{v}^n)$$

$$\Rightarrow N(\vec{v}^n) = \vec{v}^n + \Delta t \vec{F}(\vec{v}^n, t^n)$$

$$\begin{aligned} \Rightarrow N(\vec{v}^n) - \vec{v}^{n+1} &= \vec{v}^n + \Delta t \vec{F}(\vec{v}^n, t^n) - \vec{v}^{n+1} \\ &= \vec{v}^n + \Delta t \dot{\vec{v}}(t^n) - \vec{v}^{n+1} \\ &= \vec{v}^n + \Delta t \dot{\vec{v}}(t^n) - [\vec{v}^n + \Delta t \dot{\vec{v}}(t^n) + \frac{1}{2} \Delta t^2 \ddot{\vec{v}}(t^n) + O(\Delta t^3)] \\ &= -\frac{1}{2} \Delta t^2 \ddot{\vec{v}}(t^n) + O(\Delta t^3) \end{aligned}$$

$\Rightarrow \|N(\vec{v}^n) - \vec{v}^{n+1}\| \leq c \Delta t^2$  as  $\Delta t \rightarrow 0$   
 $\Rightarrow p=1$  For Forward Euler.  
 $\Rightarrow$  local accuracy of F.E. is 1<sup>st</sup> ( $p=1$ ) order  
 $\Rightarrow$  Expect global order of F.E. is 1<sup>st</sup> order.

A scheme which we have not yet studied is trapezoidal integration:

$$\frac{\vec{v}^{n+1} - \vec{v}^n}{\Delta t} = \frac{1}{2} [\vec{F}(\vec{v}^n, t^n) + \vec{F}(\vec{v}^{n+1}, t^{n+1})]$$

This is an implicit scheme since  $\vec{f}^{n+1}$  is involved. Let's look at its local accuracy:

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### Trapezoidal Local Accuracy

$$N(\vec{v}^n) = \vec{v}^n + \frac{\Delta t}{2} [\vec{f}(\vec{v}^n, t^n) + \vec{f}(\vec{v}^{n+1}, t^{n+1})]$$

$$\Rightarrow N(\vec{v}^n) - \vec{v}^{n+1} = \vec{v}^n + \frac{\Delta t}{2} [\vec{f}(\vec{v}^n, t^n) + \vec{f}(\vec{v}^{n+1}, t^{n+1})] - \vec{v}^{n+1}$$

$$= \vec{v}^n - \vec{v}^{n+1} + \frac{\Delta t}{2} [\dot{\vec{v}}(t^n) + \dot{\vec{v}}(t^{n+1})]$$

Now, Taylor series expand  $\vec{v}^{n+1} \approx \dot{\vec{v}}(t^{n+1})$ :

$$\vec{v}^{n+1} = \vec{v}^n + \Delta t \dot{\vec{v}}(t^n) + \frac{1}{2} \Delta t^2 \ddot{\vec{v}}(t^n) + \frac{1}{6} \Delta t^3 \dddot{\vec{v}}(t^n) + O(\Delta t^4)$$

$$\dot{\vec{v}}^{n+1} = \dot{\vec{v}}^n + \Delta t \ddot{\vec{v}}^n + \frac{1}{2} \Delta t^2 \dddot{\vec{v}}^n + \frac{1}{6} \Delta t^3 \overset{\cdot\cdot\cdot}{\vec{v}}^n + O(\Delta t^4)$$

$$\begin{aligned} \Rightarrow N(\vec{v}^n) - \vec{v}^{n+1} &= -\Delta t \dot{\vec{v}}^n - \frac{1}{2} \Delta t^2 \ddot{\vec{v}}^n - \frac{1}{6} \Delta t^3 \dddot{\vec{v}}^n + O(\Delta t^4) \\ &\quad + \frac{\Delta t}{2} \left[ \dot{\vec{v}}^n + \dot{\vec{v}}^n + \Delta t \ddot{\vec{v}}^n + \frac{1}{2} \Delta t^2 \ddot{\vec{v}}^n + O(\Delta t^3) \right] \\ &= \frac{1}{2} \Delta t^3 \dddot{\vec{v}}^n + O(\Delta t^4) \end{aligned}$$

$$\Rightarrow \left\| N(\vec{v}^n) - \vec{v}^{n+1} \right\| \leq C \Delta t^3 \Rightarrow p=2!$$

Trapezoidal has 2<sup>nd</sup> order accuracy

# Examples: Derivation of 2<sup>nd</sup> order Explicit Scheme

The Adams - Bashforth family of integrators are explicit and approximate the time derivatives as:

$$\left. \frac{d\vec{U}}{dt} \right|_{t^n} = \frac{\vec{U}^{n+1} - \vec{U}^n}{\Delta t}$$

1<sup>st</sup> AB (= Forward Euler):  $\frac{\vec{V}^{n+1} - \vec{V}^n}{\Delta t} = \vec{F}(\vec{V}^n, t^n)$

The 2<sup>nd</sup> A-B scheme has the form:

$$\frac{\vec{V}^{n+1} - \vec{V}^n}{\Delta t} = \beta_1 \vec{F}(\vec{V}^n, t^n) + \beta_2 \vec{F}(\vec{V}^{n-1}, t^{n-1})$$

where  $\beta_1, \beta_2$  are (currently) unknown constants  
We can use our local error analysis to determine  $\beta_{1,2}$ .

$$\vec{V}^{n+1} = \vec{V}^n + \Delta t \beta_1 \vec{F}^n + \Delta t \beta_2 \vec{F}^{n-1}$$

$$N(\vec{U}^n) - \vec{U}^{n+1} = \vec{U}^n + \Delta t \beta_1 \vec{F}(\vec{U}^n, t^n) + \Delta t \beta_2 \vec{F}(\vec{U}^{n-1}, t^{n-1}) - \vec{U}^{n+1}$$

Then, substitute and Taylor expand:

$$\vec{F}(\vec{U}^n, t^n) = \ddot{\vec{U}}^n$$

$$\vec{F}(\vec{U}^{n-1}, t^{n-1}) = \ddot{\vec{U}}^{n-1} = \ddot{\vec{U}}^n - \Delta t \dddot{\vec{U}}^n + \frac{1}{2} \Delta t^2 \overset{\dots}{\underset{\dots}{U}}^n - \frac{1}{6} \Delta t^3 \overset{\dots}{\underset{\dots}{U}}^n + O(\Delta t^4)$$

$$\vec{v}^{n+1} = \vec{v}^n + \Delta t \dot{\vec{v}}^n + \frac{1}{2} \Delta t^2 \ddot{\vec{v}}^n + \frac{1}{6} \Delta t^3 \overset{\cdot\cdot\cdot}{\vec{v}}^n + O(\Delta t^4) \quad \boxed{9}$$

Group terms in powers of  $\Delta t$ :

$$N(\vec{v}^n) - \vec{v}^{n+1} = \Delta t \dot{\vec{v}}^n (\beta_1 + \beta_2 - 1) + \Delta t^2 \ddot{\vec{v}}^n \left(-\beta_2 - \frac{1}{2}\right) \\ + \Delta t^3 \overset{\cdot\cdot\cdot}{\vec{v}}^n \left(\frac{1}{2}\beta_2 - \frac{1}{6}\right) + O(\Delta t^4)$$

Then, by setting:  $\left\{ \begin{array}{l} \beta_2 = -\frac{1}{2} \\ \beta_1 = \frac{3}{2} \end{array} \right\}$  the leading

error term is  $O(\Delta t^3)$ . Thus, the 2<sup>nd</sup> A-B scheme is:

$$\frac{\vec{v}^{n+1} - \vec{v}^n}{\Delta t} = \frac{3}{2} \vec{F}(\vec{v}^n, t^n) - \frac{1}{2} \vec{F}(\vec{v}^{n-1}, t^{n-1})$$

2<sup>nd</sup> order A-B scheme ( $p=2$ )

Note: What happens on first iteration when  $\vec{v}^{n-1}$  is not defined?

\*\* Usually use Forward Euler on first iteration.

# Summarizing Multi-Step Methods

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## Adams-Bashforth:

These are explicit schemes.

$$p=1: \frac{\vec{v}^{n+1} - \vec{v}^n}{\Delta t} = \vec{f}(\vec{v}^n, t^n) \quad (\text{Forward Euler})$$

$$p=2: \frac{\vec{v}^{n+1} - \vec{v}^n}{\Delta t} = \frac{3}{2}\vec{f}(\vec{v}^n, t^n) - \frac{1}{2}\vec{f}(\vec{v}^{n-1}, t^{n-1})$$

## Adams-Moulton:

These are implicit schemes.

$$p=1: \frac{\vec{v}^{n+1} - \vec{v}^n}{\Delta t} = \vec{f}(\vec{v}^{n+1}, t^{n+1}) \quad (\text{Backward Euler})$$

$$p=2: \frac{\vec{v}^{n+1} - \vec{v}^n}{\Delta t} = \frac{1}{2}[\vec{f}(\vec{v}^{n+1}, t^{n+1}) + \vec{f}(\vec{v}^n, t^n)] \quad (\text{Trapezoidal})$$

Variations of both of these schemes exist for  $p > 2$  also.