

Topics:

- \* Implementation of implicit schemes
- \* Accuracy and consistency

## Implementation of Implicit Schemes

Recall our system of ODE's:

$$\frac{d\vec{v}}{dt} = \vec{f}(\vec{v}, t) \quad \vec{v}(0) = \vec{v}_0$$

An implicit numerical integration is one which requires the calculation of  $\vec{f}(\vec{v}^{n+1}, t^{n+1})$  in order to calculate  $\vec{v}^{n+1}$ . For example, our Backward Euler scheme:

$$\frac{\vec{v}^{n+1} - \vec{v}^n}{\Delta t} = \vec{f}(\vec{v}^{n+1}, t^{n+1}) \quad (1)$$

We use the term "implicit" because to find  $\vec{v}^{n+1}$ , we must solve the implicit relationship between  $\vec{v}^{n+1}$  and  $\vec{f}(\vec{v}^{n+1}, t^{n+1})$ :

$$\vec{v}^{n+1} - \Delta t \vec{f}(\vec{v}^{n+1}, t^{n+1}) = \vec{v}^n$$

By comparison, an explicit scheme allows a direct

solution of  $\vec{v}^{n+1}$  without involving  $\vec{F}(\vec{v}^{n+1}, t^{n+1})$ . 12

For example, Forward Euler:

$$\frac{\vec{v}^{n+1} - \vec{v}^n}{\Delta t} = \vec{f}(t^n, \vec{v}^n)$$
$$\Rightarrow \boxed{\vec{v}^{n+1} = \vec{v}^n + \Delta t \vec{f}(t^n, \vec{v}^n)}$$

So, implicit schemes will involve more work per iteration.

### Comparison of Implicit & Explicit Schemes for Linear Systems

Suppose our canonical problem were linear:

$$\frac{d\vec{v}}{dt} = A\vec{v} + \vec{f}(t)$$

Forward Euler (FE):

$$\vec{v}^{n+1} = \vec{v}^n + \Delta t [A\vec{v}^n + \vec{f}(t^n)]$$

Thus, assuming  $\vec{f}(t)$  is not too difficult to evaluate, the main cost of FE is the matrix-vector multiply  $A\vec{v}^n$ . If  $\vec{v}$  is an  $M$  dimensional vector, this requires:

Worstcase FE work :  $M^2$  multiplications & additions

But, in general  $A$  is quite sparse.

By sparsity, we mean that the number of non-zeros in A is small compared to the number of zeros:

Define:  $M_Z = \text{number of zero entries in } A$

$M_{NZ} = \text{number of non-zero entries in } A$

$$(M^2 = M_Z + M_{NZ})$$

A sparse matrix has  $\frac{M_{NZ}}{M_Z + M_{NZ}} \ll 1$ .

\* For matrices A arising from finite difference, finite volume, or finite element discretization of a PDE, the number of non-zeros is usually a small number per row of the matrix:

$$\Rightarrow M_{NZ} = c_{NZ} M$$

where  $c_{NZ}$  = average number of non-zeros per row of A. Often  $c_{NZ} \ll M$ .

e.g. 2-D finite element for Laplace's equation would have  $c_{NZ} \approx 5$

$\Rightarrow$  The work to multiply & add the non-zero entries of A to a vector is

$$c_{NZ} M \ll M^2!$$

Now let's look at our Backward's Euler (BE) scheme [4]

$$\frac{\vec{v}^{n+1} - \vec{v}^n}{\Delta t} = A\vec{v}^{n+1} + \vec{f}(t^{n+1})$$

$$\vec{v}^{n+1} - \Delta t A \vec{v}^{n+1} = \vec{v}^n + \Delta t \vec{f}(t^{n+1})$$

$$\Rightarrow \underbrace{(I - \Delta t A)}_{M \times M \text{ matrix}} \vec{v}^{n+1} = \underbrace{\vec{v}^n + \Delta t \vec{f}(t^{n+1})}_{\text{Unknowns}} \quad \text{RHS}$$

Thus, to find  $\vec{v}^{n+1}$ , we must solve an  $M \times M$  linear system by iteration! This can get quite expensive! Some points:

- \* In the general case, Gaussian elimination (with partial pivoting) will usually work, but this requires  $M^3$  operations.
- \* Matrices arising from finite difference, element, or volume methods can often be solved more efficiently
  - structural deformations often  $O(M)$
  - heat diffusion often  $O(M)$
  - high Reynolds number air flows typically  $O(M^2)$   
but some  $O(M)$  cases
- \* For small  $\Delta t$ ,  $I - \Delta t A \approx I$ , which is obviously easy to solve. But, we want large  $\Delta t$ 's (usually if we are looking at implicit methods).

- Question: Suppose a worst case scenario such that:
- \* Ar requires  $M^2$  operations
  - \* Solving  $(I-\Delta t A)v^{n+1} = \text{RHS}$  requires  $M^3$  operations.
  - \* Assume accuracy requires  $\Delta t \leq \frac{T_{\text{long}}}{N_{\text{long}}}$

where  $T_{\text{long}} \equiv$  longest time scale of interest in problem

$N_{\text{long}} \equiv$  # of integration step required  
for desired accuracy of numerical  
integration over  $T_{\text{long}}$

Under what conditions is the implicit scheme more efficient to integrate from  $0 \rightarrow T$ ?

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The next complication is to consider nonlinear  $\vec{f}(\vec{v}, t)$  with implicit methods.

$$\text{BEs} \quad \frac{\vec{v}^{n+1} - \vec{v}^n}{\Delta t} = \vec{f}(\vec{v}^{n+1}, t^{n+1})$$

We re-write this in a residual form:

$$\vec{R}(\vec{v}) = \vec{v} - \vec{v}^n - \Delta t \vec{f}(\vec{v}, t^{n+1})$$

And we are interested in finding  $\vec{v}$  such that  $\vec{R}(\vec{v}) = 0 \Rightarrow \vec{v} = \vec{v}^{n+1}$ .

This is finding the root of a nonlinear set of coupled equations. There are several options, but we will only consider the Newton-Raphson method applied to this problem:

To find  $\vec{v}^{n+1}$ :

① Initialize  $\vec{w}^0 = \vec{v}^n$  ← use previous iteration as starting guess for  $\vec{v}^{n+1}$   
 Initialize  $m = 0$

Loop → ② Linearize about  $\vec{w}^m$

$$\vec{R}(\vec{w}^{m+1}) = \vec{R}(\vec{w}^m + \Delta \vec{w}) = 0$$

$$\approx \vec{R}(\vec{w}^m) + \left. \frac{\partial \vec{R}}{\partial \vec{w}} \right|_{\vec{w}^m} \Delta \vec{w} = 0$$

$$\Rightarrow \boxed{\left. \frac{\partial \vec{R}}{\partial \vec{w}} \right|_{\vec{w}^m} \Delta \vec{w} = -\vec{R}(\vec{w}^m)}$$

$M \times M$  linear system of eqns based

For our BE,  $\frac{\partial \vec{R}}{\partial \vec{w}}$  is:

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$$\frac{\partial \vec{R}}{\partial \vec{w}} = \frac{\partial}{\partial \vec{w}} \left\{ \vec{v} - \vec{v}^n - \Delta t \vec{f}(\vec{w}, t^{n+1}) \right\}$$

$$\frac{\partial \vec{R}}{\partial \vec{w}} = I - \Delta t \frac{\partial \vec{f}}{\partial \vec{w}}(\vec{w}, t^{n+1})$$

$$\Rightarrow \boxed{\frac{\partial \vec{R}}{\partial \vec{w}} \Big|_{\vec{w}^m} = I - \Delta t \frac{\partial \vec{f}}{\partial \vec{w}}(\vec{w}^m, t^{n+1})}$$

③ Solve linear system for  $\Delta \vec{w}$  and update:

$$\vec{w}^{m+1} = \vec{w}^m + \Delta \vec{w}$$

④ Check for convergence:

(\*) Calculate  $\vec{R}(\vec{w}^{m+1})$ . If  $\frac{\|\vec{R}(\vec{w}^{m+1})\|}{\|\vec{w}^{m+1}\|} \leq \epsilon_{tol}$   
then converged

$$\Rightarrow \boxed{\vec{v}^{m+1} = \vec{w}^{m+1} \text{ DONE}}$$

Note:  $\epsilon_{tol}$  should be small, say  $10^{-5}$   
The norm  $\|\vec{x}\|$  of  $\vec{x}$  can be done several ways:

$$* \quad \|\vec{x}\|_2 = \sqrt{\sum_{i=1}^M x_i^2} \quad \text{perhaps most common (L2 norm)}$$

$$* \quad \|\vec{x}\|_1 = \sum_{i=1}^M |x_i| \quad (\text{L1 norm})$$

$$* \quad \|\vec{x}\|_\infty = \max_{i=1 \rightarrow M} |x_i| \quad (\text{L$\infty$ norm})$$

→ (\*) If not converged sufficiently,  $m \rightarrow m+1$   
and return to ② starting next sub-iteration

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### Comments on Newton sub-iterations:

- \* In many cases, only a few sub-iterations are required as  $\vec{v}^n$  is often a good starting guess for  $\vec{v}^{n+1}$ .
- \* To be careful, you probably want to put a limit on the total number of sub-iterations.
- \* For some really tough problems, you may also want to limit the Newton sub-iteration update:

$$\vec{w}^{m+1} = \vec{w}^m + \eta \Delta \vec{w}$$

where  $0 < \eta \leq 1$ .  $\eta = 1$  is full update.  
 $\eta$  is re-calculated at every iteration usually by comparing the magnitude of some of the  $\Delta \vec{w}$  entries to  $\vec{w}^m$ . For example, suppose:

$$\vec{w} = \begin{bmatrix} X \\ Y \\ U \\ V \end{bmatrix}$$

and you had a characteristic length  $L$  and velocity  $V_\infty$

Then, you might require

$$\Delta \vec{w} = \begin{bmatrix} \Delta X \\ \Delta Y \\ \Delta U \\ \Delta V \end{bmatrix}$$

for example

$$\begin{cases} \eta \Delta X \leq 0.1 L \\ \text{or} \\ \eta \Delta U \leq 0.01 V_\infty \end{cases}$$

For work estimates:

$$FE: \frac{T}{\Delta t_{FE}} M^2$$

$$BE: \frac{T}{\Delta t_{BE}} M^3 = \frac{T}{T_{long}} N_{long} M^3$$

But  $\Delta t_{BE}$  is bound by  $\frac{T_{long}}{N_{long}}$  for accuracy

Work BE < Work FE

$$\frac{T}{T_{long}} N_{long} M^3 < \frac{T}{\Delta t_{FE}} M^2$$

$$\frac{\Delta t_{FE}}{T_{long}} N_{long} M < 1$$

But  $\Delta t_{FE} = T_{short}$  due to instability



$$\boxed{\frac{T_{short}}{T_{long}} N_{long} M < 1}$$

need stiff problems  
(this is condition number)

higher accurate numerical scheme  
will decrease  $N_{long}$

small # of states

Note: the work ratio per iteration gives rise to the M:

$$\frac{\text{Work per iter BE}}{\text{Work per iter FE}} = \frac{M^3}{M^2} = M$$

$$\text{So if } \frac{\text{Work iter BE}}{\text{Work iter PE}} = C$$

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(say if both were  $O(M)$ ) then clearly the implicit scheme is a winner for stiff problems:

$$\frac{\text{Work BE}}{\text{Work PE}} = \frac{T_{\text{short}}}{T_{\text{long}}} N \log C$$

Underlying Message: Good linear solvers (for  $Ax=b$  type problems) are critical for numerical efficiency!