

$$\underline{M} \ddot{\underline{u}} + \underline{C} \dot{\underline{u}} + \underline{K} \underline{u} = \underline{R}(t) \quad (1)$$

\underline{u}^0 $\dot{\underline{u}}^0$ given initial conditions

For solution:

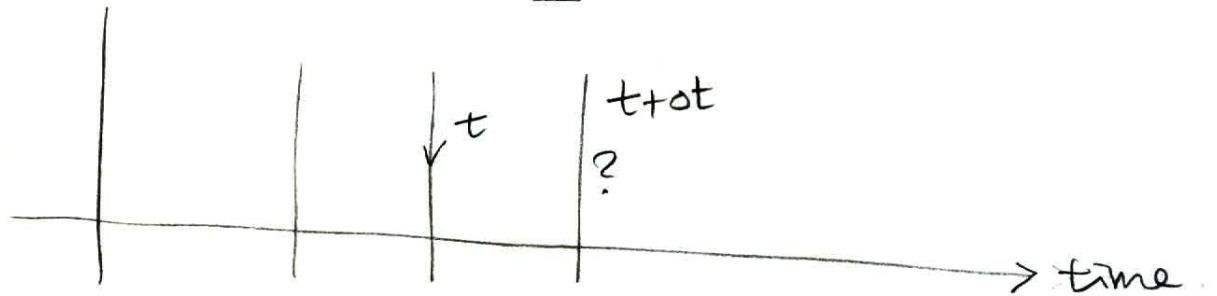
1) Explicit method, we use central difference scheme.

Properties of explicit methods

69

use $\underline{M} \underline{\ddot{u}} + \underline{c} \underline{\dot{u}} + \underline{k} \underline{u} = \underline{F}$ (a)

to obtain $\underline{u}^{t+\Delta t}$



Hence, we have conditional stability.

For CDM

$$\Delta t \leq \Delta t_{cr} = \frac{T_n}{\pi}$$

$T_n = \frac{2\pi}{\omega_n}$

We neglect damping (i.e. $\underline{c} = 0$),

We use \underline{M} = lumped mass matrix.

$${}^{t+\Delta t}u_i = {}^t\hat{R}_i \left(\frac{(\Delta t)^2}{m_{ii}} \right)$$

more on this
later!

(2) Implicit methods

Use (1) at time $t+\Delta t$, to obtain the solution at time " $t+\Delta t$ ".

We use

$$\underline{M}^{t+\Delta t} \underline{\ddot{u}} + \underline{C}^{t+\Delta t} \underline{\dot{u}} + \underline{K}^{t+\Delta t} \underline{u} = \underline{R}^{t+\Delta t} \quad (b)$$

$\uparrow \qquad \qquad \qquad \uparrow \qquad \qquad \qquad \uparrow$
 unknowns.

Newmark
method

$$\underline{\dot{u}}^{t+\Delta t} = \underline{\dot{u}}^t + \left[(1-\delta) \underline{\ddot{u}}^t + \delta \underline{\ddot{u}}^{t+\Delta t} \right] \Delta t \quad (c)$$

$$\underline{u}^{t+\Delta t} = \underline{u}^t + \underline{\dot{u}}^t \Delta t \quad (d)$$

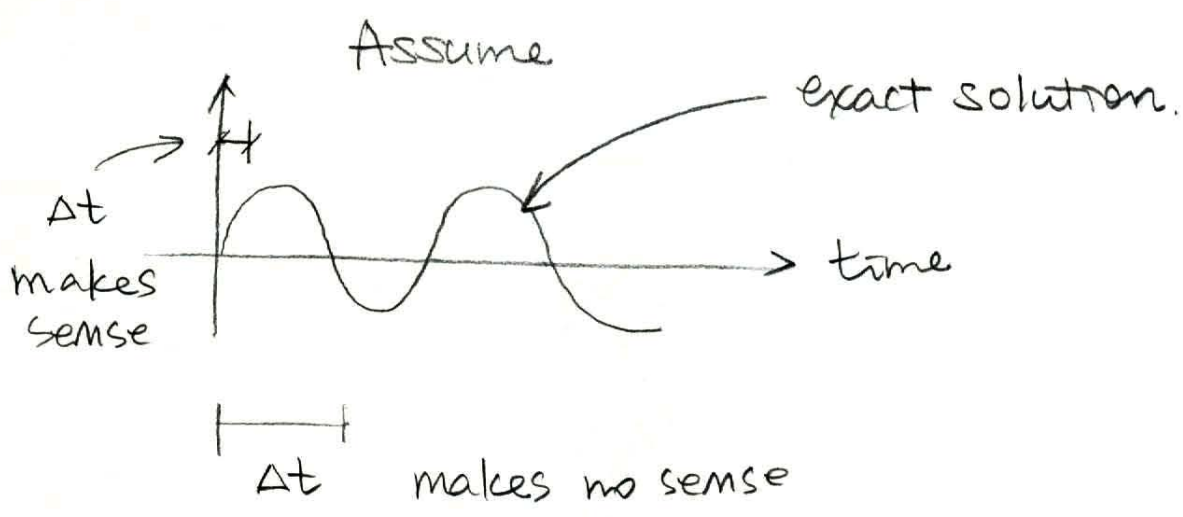
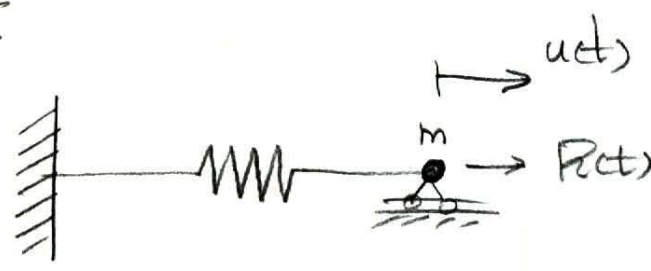
$$+ \left[\left(\frac{1}{2} - \alpha \right) \underline{\ddot{u}}^t + \alpha \underline{\ddot{u}}^{t+\Delta t} \right] (\Delta t)^2$$

Trapezoidal rule : $\delta = \frac{1}{2}$, $\alpha = \frac{1}{4}$ unconditional stability.

This means that for any Δt , the solution will not blow up.

Δt needs to be selected based on accuracy consideration.

e.g



Implicit method / trapezoidal rule

$$\hat{K}^{t+\Delta t} \underline{u} = \underline{F}^{t+\Delta t} \leftarrow \text{like static. } \star$$

$$\hat{K} = f_n \left(\underline{K}, \frac{1}{\Delta t} \underline{M}, \frac{1}{\Delta t} \underline{C} \right)$$

$\Delta t \downarrow$ effect \uparrow
 $\Delta t \uparrow$ effect \downarrow

Mode superposition, transform eq (1) into a form that is "easier to integrate."

Mode superposition

$$(A) \quad \underbrace{\underline{u}(t)}_{n \times 1} = \underbrace{\underline{P}}_{n \times n} \underline{X}(t) ; \quad \underline{P} = \text{indep. of time.}$$

$$(B) \quad \underbrace{\underline{P}^T \underline{M} \underline{P}} + \underbrace{\underline{P}^T \underline{C} \underline{P}} + \underbrace{\underline{P}^T \underline{K} \underline{P}}_{\text{Symm } \forall \underline{P}} = \underline{R}$$

We want a "P" that diagonalizes $\underline{P}^T \underline{K} \underline{P}$

$$\underline{P}^T \underline{M} \underline{P}.$$

Consider the free-vibration eqns

$$(c) \quad \underline{M} \ddot{\underline{u}} + \underline{K} \underline{u} = \underline{0}$$

Assume solution

$$\underline{u} = \underline{\phi} \sin \omega(t - t_0)$$

then substitute into (c)

$$-\underline{M} \underline{\phi} \omega^2 \sin \omega(t - t_0) + \underline{K} \underline{\phi} \sin \omega(t - t_0) = 0$$

$$\underline{K} \underline{\phi} = \omega^2 \underline{M} \underline{\phi} \quad (D)$$

For a solution to exist (D) has to be satisfied

There are n solutions to (D)

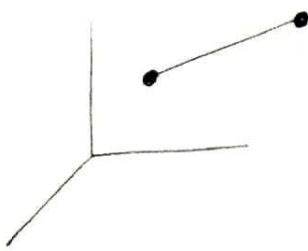
$$\underline{K} \underline{\phi}_i = \omega_i^2 \underline{M} \underline{\phi}_i \quad i=1, 2, \dots, n$$

$(\underline{\phi}_i, \omega_i^2)$ eigenpair

$$0 \leq \omega_1^2 \leq \omega_2^2 \leq \dots \leq \omega_n^2$$

$\phi_1 \quad \phi_2 \quad \dots \quad \phi_n$

eg. truss



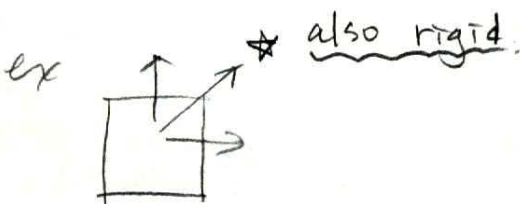
$$6 - 1 = 5 //$$

We can prove

δ_{ij} = Kronecker delta

(i) $\underline{\phi}_i^T \underline{K} \underline{\phi}_j = \omega_i^2 \delta_{ij}$

(ii) $\underline{\phi}_i^T \underline{M} \underline{\phi}_j = \delta_{ij}$



multiple same frequencies

For distinct eig values

(74)

$$\text{(eg. } \omega_{\bar{n}-1}^{\vec{\lambda}} < \omega_{\bar{n}}^{\vec{\lambda}} < \omega_{\bar{n}+1}^{\vec{\lambda}} \text{)}$$

$\underline{\phi}_{\bar{n}}$ is unique

For eigenvalues of multiplicity m ,
the eigen space of dimension m is unique
and we can always choose m eigen
vectors that satisfy (i) & (ii)

take $P = [\underline{\phi}_1, \underline{\phi}_2, \dots, \underline{\phi}_n]$